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# Stochastic aspects of oscillatory Ca<sup>2+</sup> dynamics in hepatocytes

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#### Abstract

Signal-induced  $Ca^{2+}$  oscillations have been observed in many cell types and play a primary role in cell physiology. Although it is the regular character of these oscillations that first catches the attention, a closer look at time series of  $Ca^{2+}$  increases reveals that the fluctuations on the period during individual spike trains are far from being negligible. Here, we perform a statistical analysis of the regularity of  $Ca^{2+}$  oscillations in noradrenaline-stimulated hepatocytes and find that the coefficient of variation lies between 10 and 15%. Stochastic simulations based on the Gillespie's algorithm and considering realistic numbers of  $Ca^{2+}$  ions and InsP<sub>3</sub> receptors account for this variability if the receptors are assumed to be grouped in clusters of a few tens of channels. Given the relatively small number of clusters (~200), the model predicts the existence of repetitive spikes induced by fluctuations (stochastic resonance). Oscillations of such type are found in hepatocytes at sub-threshold concentrations of noradrenaline. We next predict with the model that the isoforms of the InsP<sub>3</sub> receptor can affect the variability of the oscillations. In contrast, possible accompanying InsP<sub>3</sub> oscillations have no impact on the robustness of signal-induced repetitive  $Ca^{2+}$  spikes.

### **INTRODUCTION**

Signal-induced  $Ca^{2+}$  oscillations are observed in cells of various types and are known to play a primary role in transducing the external signal into the appropriate physiological response (1). Given the large number of physiological processes that are controlled by InsP<sub>3</sub>-induced  $Ca^{2+}$  increases, these signals are highly organized in time and space to ensure reliability and specificity (2-4). To mention only a few examples, oscillations and waves have indeed been observed in cells as various as eggs (5), cardiac myocytes (6), astrocytes (7) or plant cells (8). The nature and intensity of the external signal is encoded in the frequency, amplitude or waveform of  $Ca^{2+}$  oscillations (9,10). Given the widespread relevance of frequency encoding in  $Ca^{2+}$  dynamics, the issue of the regularity of oscillations, as that observed for many other biological rhythms (11), is of great conceptual and physiological interest.

Experimental and theoretical studies however show that randomness plays a key role in  $Ca^{2+}$  dynamics. At low stimulation levels,  $Ca^{2+}$  increases of small amplitude (~100 nM), duration (~100 ms) and spatial extent (1-3 µm) have been observed, especially in HeLa cells (12) and *Xenopus* oocytes (13). Given the small numbers of InsP<sub>3</sub>Rs and Ca<sup>2+</sup> ions taking part in these elementary events, the so-called Ca<sup>2+</sup> 'blips' and 'puffs' occur randomly. However, their frequency, amplitude and duration increase with InsP<sub>3</sub> concentration. For supra-threshold levels of stimulation, these Ca<sup>2+</sup> increases do not remain localized but propagate as intracellular waves. At the same time, Ca<sup>2+</sup> signals repeat regularly in time and can thus be called 'oscillations'. Thus, once the number of InsP<sub>3</sub>Rs that have bound InsP<sub>3</sub> and are thus susceptible to release Ca<sup>2+</sup> becomes sufficient, their Ca<sup>2+</sup> releasing activity become both coherent and periodic (12-14).

Given that both a stochastic and deterministic regime can be observed in the same cell for different levels of stimuli, one can intuitively expect that the frontier between both regimes is not clear cut and that some 'intermediate behaviour' can sometimes be observed. In particular, one can wonder to which extent global  $Ca^{2+}$  oscillations are really periodic, or if there is a significant random variation in the period during a spike train in a given cell at a fixed concentration of stimulus. Nonlinear time series analysis methods applied to experimental  $Ca^{2+}$  traces in hepatocytes suggest that stochasticity is an important factor in the dynamics of intracellular  $Ca^{2+}$  oscillations (15). This question can also be raised from the consideration that the number of membrane receptors, ion channels and  $Ca^{2+}$  ions in some membrane organelles in the cell can be very low. Following this examination, a number of modelling studies are built on stochastic simulations of  $Ca^{2+}$  dynamics (16-22). These studies have led to interesting conclusions as to the various sensitivities of the different oscillatory regimes to the low numbers of particles, or the possible predominance of the stochastic regime in the generation of apparently regular  $Ca^{2+}$  oscillations.

In the present study, we follow a combined experimental and computational approach to investigate the possible impact of stochasticity on  $Ca^{2+}$  oscillations in hepatocytes. Our analysis is based on an estimation of the regularity of  $Ca^{2+}$  oscillations in this cell type, in response to stimulation by noradrenaline. Stochastic simulations of  $Ca^{2+}$  oscillations using the Gillespie's algorithm are then developed to seek in which conditions one can recover the experimentally measured standard deviations on the period, when taking into account realistic concentrations of  $Ca^{2+}$  ions and InsP<sub>3</sub> receptors. Comparison between these simulations of the spatially averaged intracellular  $Ca^{2+}$  dynamics and experiments in the near vicinity of the bifurcation point provides convincing evidence that hepatocytes display an oscillatory regime strongly affected by internal fluctuations, due to the low number of clusters of InsP<sub>3</sub>Rs. The

impact of other factors as the  $InsP_3$  receptor isoforms or the possible accompanying  $InsP_3$  oscillations on the regularity of the  $Ca^{2+}$  spikes is also adressed.

## MATERIALS AND METHODS

## **Experiments**

#### **Materials**

Dulbecco's modified Eagle's medium and Williams'medium were from Life Technology (Invitrogen, France), Collagenase A from Boehringer (Roche Diagnostics, France). Other chemicals were purchased from Sigma (Sigma-Genosys Ltd, Sigma-Aldrich Chimie, France).

#### Hepatocyte preparation

Hepatocytes were prepared from fed female wistar rats by limited collagenase digestion of rat liver, as previously described (23). Experiments were conducted according to the CEE directives for animal experimentation (decree 2001-131; "J.O."06/02/01). After isolation, rat hepatocytes were maintained (2'106 cells/ml) at 4°C in Williams medium E supplemented with 10% fetal calf serum, penicillin (200,000U/ml) and streptomycin (100mg/ml). Cell viability, assessed by trypan blue exclusion, remained >96%, during 4-5h.

### HEK293 Cell Culture

HEK293 cells were cultured in Dulbecco's modified Eagle's medium supplemented with 10% fetal bovine serum, 100 units/ml penicillin, 100  $\mu$ g/ml streptomycin, and 2 mM glutamine. Cells were grown in an incubator at 37 °C with humidified 5% CO<sub>2</sub> and 95% air.

# Cellular Ca<sup>2+</sup> imaging

HEK293 cells grown for 48 h on 35-mm glass coverslips were loaded with 3  $\mu$ M Fura2-AM at 20°C for 30 min in saline solution containing 20mM HEPES, 116mM NaCl, 5.4mM. KCl, 1.8mM CaCl<sub>2</sub>, 0.8mM MgCl<sub>2</sub>, 0.96mM NaH<sub>2</sub>PO<sub>4</sub>, 5mM NaHCO<sub>3</sub> and 1g/L glucose (pH 7.4). Cells were then washed twice, and kept in the dark at 20 °C for at least 20 min. Determination of calcium changes in hepatocytes was as previously described (23). Briefly, hepatocytes were plated onto glass coverslips coated with type I collagen and loaded with 3 $\mu$ M Fura2-AM in modified Williams' medium, for 40min, (37°C, 5% CO<sub>2</sub>). After washing, the coverslips were transferred into a perfusion chamber placed on the stage of a Zeiss inverted microscope (Axiovert 35). Calcium imaging was performed as described previously (23). Fluorescence images were collected by a CCD camera (Princeton, USA), digitized and integrated in real time by an image processor (Metafluor, Princeton, USA).

### Computation

The numerical procedure used to simulate  $Ca^{2+}$  oscillations is based on stochastic simulations of the transitions between the various states of the InsP<sub>3</sub>-sensitive  $Ca^{2+}$  channel. These states and the possible transitions between them are schematized in Fig. 1. This model has been previously used to simulate  $Ca^{2+}$  blips and puffs (24). Many other models taking into account the positive feedback exerted by  $Ca^{2+}$  on its own release have been proposed (25). The regulation of the receptor by InsP<sub>3</sub> is not modelled dynamically: InsP<sub>3</sub> binding and unbinding indeed occurs on a faster time scale than other processes, and can thus be assumed to be always at quasi-equilibrium. Thus, the number of open InsP<sub>3</sub> receptors/Ca<sup>2+</sup> channels is given by :

$$R_{20} \frac{\left[InsP_3\right]}{K_D + \left[InsP_3\right]} \tag{1}$$

where  $R_{20}$  represents the number of InsP<sub>3</sub>Rs having 2 Ca<sup>2+</sup> ions bound at their activation site and no Ca<sup>2+</sup> ion bound at their inhibitory site (see Fig. 1) and K<sub>D</sub> the half saturation constant of InsP<sub>3</sub> for its receptor. Ca<sup>2+</sup> fluxes across the ER membrane are also modelled stochastically. If v<sub>1</sub> is the total Ca<sup>2+</sup> flux through the InsP<sub>3</sub>Rs when they are all in an open state, the corresponding deterministic evolution equation for cytosolic Ca<sup>2+</sup> concentration (C) is :

$$\frac{dC}{dt} = \beta \left[ J_{rel} - J_{serca} \right] \text{ with } J_{rel} = v_1 \frac{R_{20}}{R_T} \frac{\left[ InsP_3 \right]}{K_D + \left[ InsP_3 \right]} + v_2 \text{ and } J_{serca} = v_3 \frac{C^2}{K_3^2 + C^2}$$
(2)

where  $\beta$  stands for the effective buffering capacity of the cytoplasm and R<sub>T</sub> the total number of InsP<sub>3</sub> receptors in the cell. Such treatment of buffering relies on the assumptions that cytosolic Ca<sup>2+</sup> is buffered rapidly and that the fraction of buffered Ca<sup>2+</sup> remains constant (26). Parameter v<sub>2</sub> represents an unregulated Ca<sup>2+</sup> leak from the endoplasmic reticulum. As the concentration of Ca<sup>2+</sup> in the endoplasmic reticulum (ER) is much higher than in the cytosol, it is assumed to remain constant. Pumping from the cytosol into the ER is modelled by a Hill function (Hill coefficient equal to 2), with maximal velocity and half saturation constant represented by v<sub>3</sub> and K<sub>3</sub>, respectively.

We use the Gillespie's algorithm (27), which calculates trajectories governed by the chemical master equation and has been abundantly described elsewhere (16, 20, 28). This method of the Monte-Carlo type associates a probability to each kinetic transition considered in the reaction scheme. This probability depends on the specific stochastic reaction rate and of a combinatorial term that depends on the stoechiometry of the reaction. At each time step, the algorithm randomly determines the reaction that takes place according to its relative probability, as well as the time interval to the next reaction step. In the case of the model for  $Ca^{2+}$  oscillations, these transitions are schematized in Table 1, together with the manner by which the numbers of molecules of the different species are updated at each time step. Parameter values (Table 2) are such that, at steady state, the model reproduces the well known bell-shaped curve for the open probability of the InsP<sub>3</sub>R (29,30) as shown in Fig. 2.

### RESULTS

# Statistical analysis of Ca<sup>2+</sup> oscillations in hepatocytes

Time series of oscillations of  $Ca^{2+}$ -induced fluorescence from hepatocytes stimulated by noradrenaline were analysed. These data were obtained as explained above (see Materials and Methods). We did not consider the traces showing less than 10  $Ca^{2+}$  peaks. Typical  $Ca^{2+}$ oscillations in hepatocytes stimulated by 0.1  $\mu$ M and 1  $\mu$ M noradrenaline (Nor) are shown in Fig. 3A and 3B, respectively. For each time series, we have independently computed the mean period of  $Ca^{2+}$  oscillations and the standard deviation around this mean value. We then pooled the data obtained for cells stimulated by 0.1  $\mu$ M and 1  $\mu$ M Nor, respectively. Fig. 4 shows the resulting histograms of the coefficients of variation, CV (i.e. the standard deviation divided by the mean and expressed in %) that have been evaluated. These coefficients largely vary from one cell to the other, with an average value of the order of 13% for the lower Nor concentration and 11% for the higher concentration. The average periods of oscillations are equal to 45.5 ± 5.9 s and 26.0 ± 2.9 s for Nor 0.1  $\mu$ M and 1  $\mu$ M, respectively. These values (the average interspike interval and the coefficient of variation) were very similar for all cell preparations (n=9).

# Stochastic simulations of Ca<sup>2+</sup> oscillations in hepatocytes

In deterministic simulations and for the range of parameters listed in Table 2, sustained oscillations in Ca<sup>2+</sup> concentration occur for a finite range of InsP<sub>3</sub> concentrations (from  $0.09 \,\mu\text{M}$  to  $5.10 \,\mu\text{M}$ ). Such simulations adequately describe the real Ca<sup>2+</sup> dynamics if the number of Ca<sup>2+</sup> ions and InsP<sub>3</sub>Rs are large enough, so that internal fluctuations can be neglected. The real number of free  $Ca^{2+}$  ions can easily be evaluated from the intracellular concentration and the cytoplasmic volume of an hepatocyte. Thus, assuming a volume of  $5 \ 10^{-13}$  l for the cytoplasm of an hepatocyte (approximated as a sphere whose diameter equals 10 $\mu$ m), a basal concentration of 100 nM in the cytoplasm corresponds to ~30 000 Ca<sup>2+</sup> ions. At the top of a spike, this number is increased by a factor of the order of 10. In the approach proposed by Gillespie (27), a parameter denoted  $\Omega$  permits the modulation of the number of molecules present in the system. A value of  $\Omega$  equal to 3 10<sup>5</sup> allows us to approximatively get the appropriate numbers of free  $Ca^{2+}$  ions in our stochastic simulations (see Fig. 5). The number of InsP<sub>3</sub> receptors is more difficult to evaluate. In guinea pig hepatocytes, the amount of monomeric InsP<sub>3</sub> receptors has been estimated to  $\sim$ 190 fmol per mg of protein (31), which amounts to a density of tetrameric receptors equal to  $1.31 \ 10^{16}$  per liter. This density is about 100 times higher than in Xenopus oocytes and 200 times lower than in Purkinje cells of the cerebellum (32). In the volume of a typical hepatocyte, the density reported by Spät *et al.* (31) corresponds to about 6000 InsP<sub>3</sub>Rs. This number will determine the value of R<sub>T</sub> (total number of InsP<sub>3</sub>Rs or InsP<sub>3</sub>Rs clusters) that has to be introduced in the simulations.

Typical oscillations obtained by stochastic simulations of the model are shown in Fig. 5A. Taking into account a number of InsP<sub>3</sub>Rs in the range of what has been estimated above (5400 receptors), these oscillations are very regular. In these conditions indeed, the coefficient of variation of the period is equal to 2.4%, a value that is smaller than any value measured experimentally. Even for other values for the kinetic parameters used in the simulations, the coefficient of variation never exceeds 5%. Based on previous studies suggesting that  $InsP_3R_5$  are clustered in different cell types (24, 33), we then supposed that the channels are arranged in groups of 25. In the simulations, a cluster of n receptors is modelled as 1 channel with a n times larger conductance than an isolated receptor; indeed, if the receptors inside the cluster are assumed to be in close contact, the Ca<sup>2+</sup> concentration in their vicinity is the same and, thus, in first approximation, they all open and close simultaneously. Fig. 5B clearly shows that the regularity much decreases in these conditions; the coefficient of variation on the period is then equal to 13.9%, which is of the order of experimental estimations (see also Fig. 3A). As the number of clusters in this case is rather low (5400/25 = 216), the variability on the period much differs with the number of clusters in this range. This is shown in Fig. 6, where the total conductance of all InsP<sub>3</sub>Rs is kept constant. For large numbers of small groups of InsP<sub>3</sub>Rs (with a limit situation of 5400 single channels), the coefficient of variation becomes very small as one recovers the nearly deterministic case shown in Fig. 5A (see also point (c) in Fig. 6). Interestingly, the coefficient of variation does not vanish, even if one increases the number of clusters up to a value that corresponds to an unrealistically high number of  $InsP_3Rs$  (not shown). The remaining ~1.5% of variability can be ascribed to the relatively low number of  $Ca^{2+}$  ions. At the other extremity, i.e. for a very small number of large clusters, the oscillations become unrealistically irregular (see inset corresponding to point (a) in Fig.6).

As expected, variability also decreases with the concentration of  $InsP_3$ . A decrease in  $InsP_3$  concentration indeed corresponds to a lower number of channels taking an active part in  $Ca^{2+}$  release. Below the deterministic oscillatory regime, sub-threshold  $InsP_3$  concentrations lead to widely spaced, very irregular spikes; this is visible in the first part of Fig. 7A where

the level of  $InsP_3$  is just below the bifurcation point delimiting the oscillatory domain, and would thus correspond to a stable steady state in the deterministic regime. However, as this state is excitable, when fluctuations are allowed to occur, they can sometimes pass the threshold for excitability and generate a whole  $Ca^{2+}$  spike. This stochastic resonance phenomenon, which can only be obtained with a low number of  $Ca^{2+}$ -releasing channels, has also been observed in other models for  $Ca^{2+}$  oscillations (20, 33, 34). Such noise-induced  $Ca^{2+}$ oscillations are characterized by large standard deviations on the mean interspike interval (35% in the case of Fig. 7A). To compare, just on the other side of the bifurcation point ([InsP<sub>3</sub>] =  $0.095 \,\mu$ M), oscillations are much more regular (17% in the case of the second part of Fig. 7A). Such a steep change in the regularity of  $Ca^{2+}$  oscillations can sometimes be observed in hepatocytes, when the level of noradrenaline is finely adjusted at the border of the oscillatory domain (Fig. 7B). A sub-threshold concentration of noradrenaline (0.03 µM) induces noisy and irregular  $Ca^{2+}$  oscillations (CV=31%). A slight increase of the nodadrenaline concentration (0.05  $\mu$ M) then provokes a significant increase in the regularity of oscillations (CV = 12%), which becomes qualitatively similar to the repetitive spikes observed in the whole oscillatory domain (see Fig. 4). The results shown in Fig. 7 reveal two important dynamical characteristics of  $Ca^{2+}$  oscillations in hepatocytes : first that the number of clusters of InsP<sub>3</sub> receptors is low, allowing the existence of noise-induced Ca<sup>2+</sup> oscillations, and, second, that for most stimulation levels, oscillations occur in a deterministically oscillatory regime, where the coefficient of variation, CV, is much smaller than for noiseinduced  $Ca^{2+}$  oscillations.

# Other factors affecting the regularity of Ca<sup>2+</sup> oscillations

Observations of  $Ca^{2+}$  oscillations in different cell types suggest that their regularity varies from one cell type to the other (35). Many factors can be involved in this effect, such as the size of the cell and the receptor density. The respective amounts of the InsP<sub>3</sub>R isoforms could also play a role in this respect. Three isoforms of this channel, differing in their regulatory properties by  $Ca^{2+}$  and InsP<sub>3</sub>, have indeed been identified. These are co-expressed within cells, but their respective levels of expression are largely tissue- and development-specific (36). The amounts of each isoform of a given subtype have been modified by genetic manipulations in DT40, HeLa and COS-7 cells, leading to the idea that there is a close correlation between the types of InsP<sub>3</sub>Rs present in a cell and the existence, characteristics and regularity of  $Ca^{2+}$  oscillations (37, 38).

In a previous study based on a deterministic approach (39), we have shown that slight modifications in the regulatory properties of the InsP<sub>3</sub>Rs can lead to significantly different oscillatory properties when their respective densities are varied. Here, we perform stochastic simulations to test whether the robustness of  $Ca^{2+}$  oscillations can also be affected by the receptor subtype. Thus, we change the parameters values characterizing the InsP<sub>3</sub>R dynamics in order to change the shape of the bell-shaped curve showing the channel opening probability (Fig. 2). Qualitatively, the changes performed correspond to a channel that is less sensitive to  $Ca^{2+}$  changes (wider bell-shaped curve) and that is activated by slightly higher  $Ca^{2+}$  concentrations. In reality, this subtype is associated with a dependence of the level of  $Ca^{2+}$  corresponding to the highest opening probability on the InsP<sub>3</sub> level ('shift' of the bell-shaped curve with the InsP<sub>3</sub> concentration), but as the present simulations are performed with a constant InsP<sub>3</sub> concentration, it is only reflected by different kinetic parameters leading to a wider bell-shaped curve (rate of  $Ca^{2+}$  binding to and unbinding from the inhibitory binding site, see Table 2).

At low stimulation levels, Ca<sup>2+</sup> oscillations generated by the InsP<sub>3</sub>R characterized by the wider bell-shaped curve show approximately the same coefficient of variation as those

obtained before (Fig. 8). Interestingly, the coefficient of variation rapidly increases with  $InsP_3$  concentration. Intuitively, one can understand that, as  $Ca^{2+}$ -induced inhibition is less efficient (wider bell-shaped curve), these  $InsP_3R$  can sometimes remain open longer, leading to prolonged  $Ca^{2+}$  peaks as those shown in the inset of Fig. 8. This in turn induces a higher variability of the period. From a different point of view, as a widening of the bell-shaped curve is associated with a decrease in the range of  $InsP_3$  concentrations leading to oscillations, fluctuations occasionally push the system out of the oscillatory regime, thus leading to a incident increase of the spike width and thereby of the interspike interval. Interestingly, this also induces an increase of the coefficient variation with the level of  $InsP_3$ , while the opposite was obtained for the other parameter values, as well as observed in hepatocytes.

In a hypothetical way, one can assume that the wider bell-shaped curve is approaching the type 1 InsP<sub>3</sub>R, while the original parameters would qualitatively more closely correspond to the InsP<sub>3</sub>R2, which is predominantly expressed in hepatocytes. In this framework, indirect confirmation of these higher values for the coefficient of variation when widening the bellshaped curve is given by the analysis of  $Ca^{2+}$  oscillations in HEK (human embryonic kidney) cells, a cell line that expresses only type 1 and type 3 receptors. An analysis of time series of  $Ca^{2+}$  oscillations in HEK cells stimulated either by ATP or carbachol indeed indicates an average coefficient of variation equal to 31% (n=30). An example of such time series is given in Fig. 9, where a broadening of the spikes similar to what is seen in the simulations can be observed. A similar value for the coefficient of variation has been estimated from the analysis of the interspike intervals in HEK cells stimulated carbachol (35).

Finally, we have used our stochastic model to investigate the effect of potential InsP<sub>3</sub> variations accompanying Ca<sup>2+</sup> oscillations on the regularity of the interspike intervals. Observations performed in epithelial (40) or CHO (41) cells indeed suggest that Ca<sup>2+</sup> and InsP<sub>3</sub> could oscillate in synchrony. It is plausible that both oscillatory mechanisms (InsP<sub>3</sub>-driven or not) could coexist in different cell types (42). InsP<sub>3</sub> oscillations can result from regulation by Ca<sup>2+</sup> of either InsP<sub>3</sub> synthesis or InsP<sub>3</sub> degradation; PLC, the enzyme responsible for InsP<sub>3</sub> synthesis, indeed requires Ca<sup>2+</sup> for its activity, and, for some isoforms, stimulation of PLC activity by Ca<sup>2+</sup> occurs in the same concentration range as Ca<sup>2+</sup> calmodulin complex (43). Although this regulation is thought to play a minor role in the existence and characteristics of Ca<sup>2+</sup> oscillations in hepatocytes (23), one might expect that it could lead to a stabilization of the period of oscillations.

We have thus included additional steps in the Gillespie's algorithm to model InsP<sub>3</sub> synthesis and degradation. The kinetic expressions used are of the Michaelis type (see legend to Fig. 10). The number of InsP<sub>3</sub> molecules (of the order of  $10^5$ ) was chosen to realistically model an InsP<sub>3</sub> concentration of the order of 1  $\mu$ M. When Ca<sup>2+</sup> stimulates InsP<sub>3</sub> synthesis, we found no significant effect of InsP<sub>3</sub> dynamics on the robustness of Ca<sup>2+</sup> oscillations (see Fig. 10): the coefficient of variation is in the range 10-15% in the whole oscillatory domain. Other simulations where InsP<sub>3</sub> oscillations are provoked by the Ca<sup>2+</sup>-induced activation of InsP<sub>3</sub> transformation into InsP<sub>4</sub> lead to the same conclusion. As in the simulations shown above, most internal fluctuations are indeed due to the species that is present in the lowest amount, i.e. the number of clusters of InsP<sub>3</sub> receptors.

### Discussion

In this study, we have shown that  $Ca^{2+}$  spikes in hepatocytes are intrinsically irregular, as the spikes are paced with a precision that does not exceed 85%. Stochastic simulations incorporating realistic numbers of  $Ca^{2+}$  ions and InsP<sub>3</sub>Rs argue that this irregularity can be ascribed to the gathering of InsP<sub>3</sub>Rs in groups of a few tens of channels. The concept of

clustering of Ca<sup>2+</sup> channels is well-known to be necessary to explain the characterisics of smaller-scale  $Ca^{2+}$  increases known as ' $Ca^{2+}$  puffs' (17, 24, 33). The present analysis indicates that such a clustering is also necessary to account for the intrinsic irregularity of the repetitive, global  $Ca^{2+}$  spikes. Interestingly, it has also been shown that  $Ca^{2+}$  signalling capability of the cell is modified with the distribution of the  $Ca^{2+}$  release channels; channel clustering can indeed enhance the cell's capability to generate a large response to a weak InsP<sub>3</sub> signal (44). Taken together, this study and the results obtained by Shuai and Jung (2003) indicate that the distribution of InsP<sub>3</sub>Rs in clusters is a compromise between optimizing the sensitivity of the cell to weak stimuli and ensuring robust oscillations. Given the important role played by internal fluctuations at the low number of clusters predicted by the oscillations (~200), noiseinduced Ca<sup>2+</sup> oscillations sometimes take place in hepatocytes at sub-threshold concentrations of noradrenaline. In agreement with the limited InsP<sub>3</sub> range in which such noisy repetitive spiking can be observed in the simulations, this behavior is in fact rarely observed in experiments. In other cell types where the density of clusters is larger, such noise-induced Ca<sup>2+</sup> oscillations are not expected to occur, as the impact of internal fluctuations would be much reduced in this case.

The present results also suggest that for most stimulation levels, experimentally observed  $Ca^{2+}$  oscillations in hepatocytes correspond to an oscillatory regime; the steep increase in the coefficient of variation shown in Fig. 7 indeed corresponds to the passage from an excitable to an oscillatory regime. Studies in other cell types have however led to the opposite conclusion (17, 19). Instead,  $Ca^{2+}$  oscillations are there viewed as a sequence of random spikes (35). This dissent may find its origin in the fact that the oscillations reported in the latter study are in most cases spontaneous (i.e. not induced by the application of any hormone); these noisy oscillations may thus rely on another type of dynamics (45). It is on the other hand meaningful that different cell types use different ways to display repetitive Ca<sup>2+</sup> spiking because of structural disparities. In the simulations of  $Ca^{2+}$  waves in *Xenopus* oocytes that have led to suggest a stochastic nature to  $Ca^{2+}$  oscillations (17), the so-called 'focal sites', made of a group of nearby clusters, are responsible for the nucleation of the wave. In this framework, an entire hepatocyte can be viewed as a large focal site, as the mean average distance between clusters is of the order of 1  $\mu$ m (estimated on the basis of ~200 clusters). In this respect, it is interesting to mention that puffs have never been reported in hepatocytes. Although one cannot exclude that their observation is in fact limited by technical considerations, this may suggest that the clusters arrangement is such that the firing of any of them will automatically induce a global cytoplasmic Ca<sup>2+</sup> increase. Other cell to cell differences in the regularity of the spikes can be ascribed to the different populations of InsP<sub>3</sub>Rs isoforms, which are known to characterize the various cell types (36). Our results indeed suggest that oscillations relying on the type 2 InsP<sub>3</sub> receptor will be particularly robust, which qualitatively fits with the experimental observations peformed by Miyakawa et al. (37) in DT40 B-cells and Morel et al. (46) in myocytes.

The period of the simulated  $Ca^{2+}$  oscillations presented here is always shorter than in the experiments. This drawback may probably be avoided by taking into account other regulatory processes such as  $Ca^{2+}$  exchange with the extracellular medium (47),  $Ca^{2+}$  handling by mitochondria (48), the detailed kinetics of the hormonal receptor (18), local luminal  $Ca^{2+}$ depletion (49) or a dynamic modelling of buffering (50,51). As shown in previous studies (18), the consideration of this latter factor would slightly decrease the fluctuations in free  $Ca^{2+}$ . Although we are aware of the fact that the model is simplified with respect to all these additional controls, our present hypothesis is that such processes would not affect our conclusions as to the robustness of the oscillations, which mainly depends on the number of clusters of InsP<sub>3</sub>Rs. In the same manner, we have neglected spatial aspects in the simulations, which are necessary to account for wave propagation. Incorporation of  $Ca^{2+}$  diffusion in stochastic modelling is very time-consuming. In most cases, simplifying hypotheses have been introduced (17, 49). Another consequence of neglecting spatial aspects is that our representation of clusters where all InsP<sub>3</sub>Rs/Ca<sup>2+</sup> channels open and close together is also oversimplified. More realistic descriptions of clusters can be found elsewhere (17, 19, 21, 24, 33, 49). We plan to extend the present work to incorporate Ca<sup>2+</sup> diffusion into the present Gillespie's algorithm, as it has been done for the minimal Ca<sup>2+</sup>-induced Ca<sup>2+</sup> release model (16).

As, *in vivo*, hepatocytes are connected and coupled through gap junctions and display synchronized  $Ca^{2+}$  spikes (52), incorporation of  $Ca^{2+}$  diffusion will also allow to study the effect of intercellular coupling on the robustness of  $Ca^{2+}$  oscillations in hepatocytes. The importance of stochastic effects in modelling  $Ca^{2+}$  oscillations in connected hepatocytes has already been assessed (53). Of particular interest is the conclusion reached by these authors that fluctuations in intracellular InsP<sub>3</sub> and  $Ca^{2+}$  levels decrease the threshold level of gap junction permeability necessary to coordinate  $Ca^{2+}$  spiking between adjacent cells. *In vivo*, an additional level of complexity arises from the fact that the hormonal stimulus does not remain strictly constant, and in some instances, even follows an oscillatory pattern (11, 54). Thus, as it is nowadays largely emphasized for genetic systems (28, 55), molecular noise appears to be an important component of the oscillatory  $Ca^{2+}$  dynamics, which has to be considered for a detailed elucidation of this widespread signalling pathway.

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<b>Reaction step</b>	Probability	Changes in particle numbers
$R_{00} + C \rightarrow R_{10}$	$k_{a1^+} \ C \ R_{00} \ / \ \Omega$	$R_{00} \rightarrow R_{00} - 1; R_{10} \rightarrow R_{10} + 1; C \rightarrow C - 1$
$R_{10} \rightarrow R_{00} + C$	$k_{a1-} R_{10}$	$R_{10} \rightarrow R_{10} - 1; R_{00} \rightarrow R_{00} + 1; C \rightarrow C + 1$
$R_{10} + C \twoheadrightarrow R_{20}$	$k_{a2^+} \ C \ R_{10} \ / \ \Omega$	$R_{10} \rightarrow R_{10} - 1; R_{20} \rightarrow R_{20} + 1; C \rightarrow C - 1$
$R_{20} \rightarrow R_{10} + C$	$k_{a2-} R_{20}$	$R_{20} \rightarrow R_{20} - 1; R_{10} \rightarrow R_{10} + 1; C \rightarrow C + 1$
$R_{01} + C \rightarrow R_{11}$	$k_{a1^+} \ C \ R_{01} \ / \ \Omega$	$R_{01} \rightarrow R_{01} - 1; R_{11} \rightarrow R_{11} + 1; C \rightarrow C - 1$
$R_{11} \rightarrow R_{01} + C$	k <sub>a1-</sub> R <sub>11</sub>	$R_{11} \rightarrow R_{11} - 1; R_{01} \rightarrow R_{01} + 1; C \rightarrow C + 1$
$\mathbf{R}_{11} + \mathbf{C} \twoheadrightarrow \mathbf{R}_{21}$	$k_{a2^+} \ C \ R_{11} \ / \ \Omega$	$R_{11} \rightarrow R_{11} - 1; R_{21} \rightarrow R_{21} + 1; C \rightarrow C - 1$
$R_{21} \rightarrow R_{11} + C$	$k_{a2-} R_{21}$	$R_{21} \rightarrow R_{21} - 1; R_{11} \rightarrow R_{11} + 1; C \rightarrow C + 1$
$R_{02} + C \rightarrow R_{12}$	$k_{a1^+} \ C \ R_{02} \ / \ \Omega$	$R_{02} \rightarrow R_{02} - 1; R_{12} \rightarrow R_{12} + 1; C \rightarrow C - 1$
$R_{12} \rightarrow R_{02} + C$	$k_{a1-} R_{12}$	$R_{12} \rightarrow R_{12} - 1; R_{02} \rightarrow R_{02} + 1; C \rightarrow C + 1$
$R_{12} + C \twoheadrightarrow R_{22}$	$k_{a2^+} \ C \ R_{12} \ / \ \Omega$	$R_{12} \rightarrow R_{12} - 1; R_{22} \rightarrow R_{22} + 1; C \rightarrow C - 1$
$R_{22} \rightarrow R_{12} + C$	$k_{a2-} R_{22}$	$R_{22} \rightarrow R_{22} - 1; R_{12} \rightarrow R_{12} + 1; C \rightarrow C + 1$
$R_{00} + C \rightarrow R_{01}$	$k_{i1+} \ C \ R_{00} \ / \ \Omega$	$R_{00} \rightarrow R_{00} - 1; R_{01} \rightarrow R_{01} + 1; C \rightarrow C - 1$
$\mathbf{R}_{01} \rightarrow \mathbf{R}_{00} + \mathbf{C}$	$k_{i1-} R_{01}$	$R_{01} \rightarrow R_{01} - 1; R_{00} \rightarrow R_{00} + 1; C \rightarrow C + 1$
$R_{01} + C \twoheadrightarrow R_{02}$	$k_{i2^+} \ C \ R_{01} \ / \ \Omega$	$R_{01} \rightarrow R_{01} - 1; R_{02} \rightarrow R_{02} + 1; C \rightarrow C - 1$
$R_{02} \rightarrow R_{01} + C$	$k_{i2}$ - $R_{02}$	$R_{02} \rightarrow R_{02} - 1; R_{01} \rightarrow R_{01} + 1; C \rightarrow C + 1$
$R_{10} + C \rightarrow R_{11}$	$k_{i1+} \ C \ R_{10} \ / \ \Omega$	$R_{10} \rightarrow R_{10} - 1; R_{11} \rightarrow R_{11} + 1; C \rightarrow C - 1$
$\mathbf{R}_{11} \rightarrow \mathbf{R}_{10} + \mathbf{C}$	$k_{i1}$ $R_{11}$	$R_{11} \rightarrow R_{11} - 1; R_{10} \rightarrow R_{10} + 1; C \rightarrow C + 1$
$R_{11} + C \rightarrow R_{12}$	$k_{i2^+} C  R_{11} / \Omega$	$R_{11} \rightarrow R_{11} - 1; R_{12} \rightarrow R_{12} + 1; C \rightarrow C - 1$
$R_{12} \rightarrow R_{11} + C$	k <sub>i2-</sub> R <sub>12</sub>	$R_{12} \rightarrow R_{12} - 1; R_{11} \rightarrow R_{11} + 1; C \rightarrow C + 1$
$R_{20} + C \twoheadrightarrow R_{21}$	$k_{i1+} \ C \ R_{20} \ / \ \Omega$	$R_{20} \rightarrow R_{20} - 1; R_{21} \rightarrow R_{21} + 1; C \rightarrow C - 1$
$\mathbf{R}_{21} \rightarrow \mathbf{R}_{20} + \mathbf{C}$	k <sub>i1-</sub> R <sub>21</sub>	$R_{21} \rightarrow R_{21} - 1; R_{20} \rightarrow R_{20} + 1; C \rightarrow C + 1$
$R_{21} + C \rightarrow R_{22}$	$k_{i2^+} C \; R_{21} \; / \; \Omega$	$R_{21} \rightarrow R_{21} - 1; R_{22} \rightarrow R_{22} + 1; C \rightarrow C - 1$
$R_{22} \rightarrow R_{21} + C$	k <sub>i2-</sub> R <sub>22</sub>	$R_{22} \rightarrow R_{22} - 1; R_{21} \rightarrow R_{21} + 1; C \rightarrow C + 1$
$C_{ER} \xrightarrow{InsP_3R} C$	$V_1 \frac{R_{22}}{R_{22}} \frac{[InsP_3]}{R_{22}}$	$C \rightarrow C+1$
	$R_T  K_D + \lfloor InsP_3 \rfloor$	
$C_{ER} \xrightarrow{leak} C$	$v_2 \Omega$	$C \rightarrow C+1$
$C \rightarrow C_{ER}$	$v_3 \frac{C^2}{K_3^2 + C^2} \ \Omega$	$C \rightarrow C-1$

**Table 1**. Stochastic model for  $Ca^{2+}$  oscillations. The first column lists the sequence of reactions; the probability of each reaction to occur within an infinitesimal time interval is given in the second column. The last column indicates the changes in the number of molecules/ions taking part in the different reactions.

Parameter	Description	Value
$k_{a1+}$	$Ca^{2+}$ binding to the first activating site of the InsP <sub>3</sub> R	350 μM <sup>-1</sup> s <sup>-1</sup>
k <sub>a2+</sub>	Ca <sup>2+</sup> binding to the second activating site of the InsP <sub>3</sub> R	$20000 \mu M^{-1} s^{-1}$
k <sub>a1-</sub>	$Ca^{2+}$ dissociation from the first activating site of the InsP <sub>3</sub> R	$3000 \text{ s}^{-1}$
k <sub>a2-</sub>	Ca <sup>2+</sup> dissociation from the second activating site of the InsP <sub>3</sub> R	$30 \text{ s}^{-1}$
$k_{i1+}$	Ca <sup>2+</sup> binding to the first inhibiting site of the InsP <sub>3</sub> R	$0.5 (0.2) \mu M^{-1} s^{-1}$
$k_{i2+}$	Ca <sup>2+</sup> binding to the second inhibiting site of the InsP <sub>3</sub> R	$100 (20) \mu M^{-1} s^{-1}$
k <sub>i1-</sub>	Ca <sup>2+</sup> dissociation from the first inhibiting site of the InsP <sub>3</sub> R	$25 \text{ s}^{-1}$
k <sub>i2-</sub>	Ca <sup>2+</sup> dissociation from the second inhibiting site of the InsP <sub>3</sub> R	$0.2 \text{ s}^{-1}$
β	Ca <sup>2+</sup> buffering capacity of the cytoplasm	0.05
K <sub>D</sub>	Half saturation constant of InsP <sub>3</sub> for its receptor	0.35 µM
$\mathbf{v}_1$	Maximal rate of Ca <sup>2+</sup> release through the InsP <sub>3</sub> receptor	600 μM <sup>-1</sup> s <sup>-1</sup>
$\mathbf{v}_2$	Ca <sup>2+</sup> leak from the endoplasmic reticulum	$2 \mu M^{-1} s^{-1}$
<b>V</b> 3	Maximal rate of Ca <sup>2+</sup> pumping into the ER	$100 \ \mu M^{-1} s^{-1}$
$K_3$	Half saturation constant of Ca <sup>2+</sup> pumping into the ER	0.1 µM
R <sub>T</sub>	Total number of clusters of InsP <sub>3</sub> Rs considered in the simulations	9-5400
Ω	Parameter of the Gillespie's algorithm allowing to modulate the number of $Ca^{2+}$ ions	3 10 <sup>5</sup>

**Table 2**. Parameter values used in the Gillespie simulations of  $Ca^{2+}$  oscillations described in the 'Materials and Methods' section and in Table 1. Values in parentheses refer to the modelling of the other isoform of the InsP<sub>3</sub> receptor shown in Fig. 8. The value of parameter  $R_T$  is given in the legends for each figure.

### **Figure legends**

**Fig. 1** Schematic representation of the model used to simulate the dynamics of the InsP<sub>3</sub> receptor/ Ca<sup>2+</sup> channel. The channel exhibits different states, depending on the absence or presence of Ca<sup>2+</sup> at the binding sites: Rij refers to the state of the channel with i Ca<sup>2+</sup> ions bound at the activating site and j Ca<sup>2+</sup> ions bound at the inhibiting site. Binding of InsP<sub>3</sub> is assumed to be always at equilibrium. Cooperativity in Ca<sup>2+</sup> binding at both sites is accounted for by the fact that  $k_{a2+} \gg k_{a1+}$ ,  $k_{a2-} \ll k_{a1-}$ ,  $k_{i2+} \gg k_{i1+}$  and  $k_{i2-} \ll k_{i1-}$ .

**Fig. 2** Stationary open probability of the InsP<sub>3</sub> receptor obtained by solving the evolution equations of the model schematized in Fig.1 at steady-state. The 2 curves correspond to different isoforms of the receptor. The full line is compatible with a form of the receptor that is very sensitive to  $Ca^{2+}$  changes (possibly type 2) while the dashed line would correspond to a wider bell-shaped curve, which is slightly shifted to the right (possibly type 1). The inset corresponds to the same curves shown in linear scale. The receptor whose behavior is illustrated by the full line has been used for all figures, except for Fig. 8 where parameters values corresponding to the dashed line have been used. Parameter values are those listed in Table 2 with [InsP<sub>3</sub>] = 0.1  $\mu$ M, with the dashed line corresponding to the values indicated in parentheses. Both curves are normalized with respect to their maximum, i.e. 0.19 for the full line and 0.79 for the dashed line. See (Dupont and Combettes, 2007) for more details about the modelling of the InsP<sub>3</sub>R isoforms.

**Fig. 3** Typical  $Ca^{2+}$  oscillations in noradrenaline-stimulated hepatocytes. The concentration of noradrenaline is 0.1  $\mu$ M in panel (A) and 1  $\mu$ M in panel (B).

**Fig. 4** Histograms of the coefficients of variations calculated from experimental time-series of  $Ca^{2+}$  oscillations in hepatocytes stimulated by noradrenaline. The upper panel synthetizes the measurements performed in 68 cells, all stimulated by 0.1  $\mu$ M noradrenaline. The average coefficient of variation equals 13%, and the average period 45.5 s. For the lower panel (n=47), other cells were stimulated by 1  $\mu$ M noradrenaline. The average coefficient of variation and period equal 11% and 26.0 s, respectively. For both panels, all cells considered displayed more than 10  $Ca^{2+}$  spikes whose maximum was always larger than the average fluorescence.

**Fig. 5** Gillespie's simulations of  $Ca^{2+}$  oscillations. (A) All InsP<sub>3</sub>Rs (5400) are considered to be independently regulated by  $Ca^{2+}$ . (B) The InsP<sub>3</sub>Rs are assumed to be clustered in groups of 25 channels, thus,  $R_T = 216$ . Each group of 25 channels is modelled, in first approximation, as one 'mega-channel' as it is assumed that all channels open simultaneously, and subsequently become inhibited at the same time. The conductance of this mega-channel is 25 times the conductance of a single InsP<sub>3</sub>R. For both panels, reactions and parameters are given in Tables 1 and 2. [InsP<sub>3</sub>] = 0.2  $\mu$ M.

**Fig. 6** Relationship between the coefficient of variation and the number of clusters,  $R_T$  in the stochastic model for  $Ca^{2+}$  oscillations. The total  $Ca^{2+}$  flux through all the InsP<sub>3</sub>Rs is assumed to remains constant. The insets show oscillations on a 100s period of time, obtained with 45 (a), 216 (b) or 5400 (c) clusters. All the simulations are performed with reactions and parameters given in Tables 1 and 2, with [InsP<sub>3</sub>] = 0.2  $\mu$ M.

Fig. 7 Transition from noise-induced  $Ca^{2+}$  oscillations to deterministic  $Ca^{2+}$  oscillations in hepatocytes. Panel (A) shows the results of stochastic simulations at an InsP<sub>3</sub> concentration

just below the bifurcation point for 10 min. The level of  $InsP_3$  is then instantaneously increased by 0.01  $\mu$ M, which provokes the entry in the deterministic oscillatory regime. This change of dynamical regime is accompanied by a large decrease of the coefficient of variation. Panel (B) shows an experimental trace of fluorescence obtained in an hepatocyte at 2 very low concentrations of noradrenaline, which presumably correspond to the passage through the bifurcation point for this particular cell. Such a behaviour is rarely observed (less than 10% of the cells responding to this low concentration of noradrenaline).

**Fig. 8** Effect of changing the isoform of the InsP<sub>3</sub> receptor on the robustness of Ca<sup>2+</sup> oscillations. Shown are the mean interspike interval ( $\blacksquare$ , dashed line) and coefficient of variation ( $\blacktriangledown$ , plain line) obtained by Gillespie simulations of the model detailed in Table 1, with parameter values shown in Table 2 considering the values indicated in parentheses. As shown by the dashed curve of Fig. 2, these parameters correspond to an InsP<sub>3</sub>R isoform that is less sensitive to Ca<sup>2+</sup> changes. The dependence of the rate of inhibition of the InsP<sub>3</sub>R of type 1 on the level of InsP<sub>3</sub> is not considered in the simulations, as it does not influence the robustness of oscillations at constant InsP<sub>3</sub> concentration. The inset shows typical Ca<sup>2+</sup> oscillations obtained in the simulations. For the inset, [InsP<sub>3</sub>] = 0.12  $\mu$ M and the mean interspike interval equals 8.8s; a 100s simulation is shown.

**Fig. 9** Typical oscillations in cytosolic  $Ca^{2+}$  in HEK cells stimulated by carbachol. The indicated value of the coefficient of variation, CV, has been established on a total of 54 peaks (representing 55 min of monitoring), among which only 12 are shown here. This value is close to the average value of 31% obtained for 30 cells.

**Fig. 10** Gillespie simulations of Ca<sup>2+</sup> when including InsP<sub>3</sub> oscillations resulting from the positive feedback exerted by Ca<sup>2+</sup> on PLC activity. Transition rates and parameters are as in Fig. 5, with additional steps corresponding to the stochastic version of the following evolution equation :  $\frac{d[InsP_3]}{dt} = \gamma V_{PLC} \frac{C}{K_A + C} - V_d \frac{[InsP_3]}{K_d + [InsP_3]} - k_{NS}[InsP_3]$  with V<sub>PLC</sub> = 100 µMs<sup>-1</sup>; K<sub>A</sub> = 0.3 µM ; V<sub>d</sub> = 5 µMs<sup>-1</sup> ; K<sub>d</sub> = 0.6 µM ; k<sub>NS</sub> = 0.5 s<sup>-1</sup> and  $\gamma$  = 0.07. For these parameter values, the coefficient of variation equals 13.6%.



Figure 1



Figure 2





Figure 3





Figure 5



Figure 6



Figure 7



Figure 8



Figure 9



Figure 10